

Physics 215A QFT Fall 2026 Assignment 2

Due 12:30pm Thursday, October 8, 2026

Thanks in advance for following the submission guidelines on hw01. Please ask me by email if you have any trouble.

A comment about what to hand in: please feel free to skip straightforward steps in your final writeup to save yourself some typing.

1. **Brain-warmer on units.** Briefly show that we can convert between energy, momentum, inverse length, and inverse time by multiplying by various factors of $\hbar$ and c . Thus these dimensionful quantities are all directly comparable once we choose units where $\hbar = c = 1$. What are the length scale and time scale and mass scale associated with $1 \text{ GeV} = 10^9 \text{ eV}$? What are the units of the Newton constant of gravity, and what is its value expressed in terms of powers of GeV?
2. **The correct way to think about equations of motion.** Suppose we have an action $S[\{\phi_r\}]$ that depends on a collection of fields $\phi_r(x)$ (say they live in D spacetime dimensions). The equation of motion is the condition for stationary phase in the path integral:

$$0 = \frac{\delta S}{\delta \phi_r(x)} \quad \text{for each } r. \quad (1)$$

Consider the special but important and familiar case when

- S is local: $S[\{\phi_r\}] = \int d^D x \mathcal{L}(\phi_r(x), \partial_\mu \phi_r(x), \partial_\mu \partial_\nu \phi_r(x), \dots)$ where the Lagrange density $\mathcal{L}$ depends only on the fields at x and their derivatives at x , and
- $\mathcal{L}$ depends only on ϕ and $\partial_\mu \phi$ (*i.e.* only first derivatives, and no higher derivatives of ϕ).

Show (using only the fact that the fields at different spacetime points are independent ($\frac{\delta \phi_r(x)}{\delta \phi_s(y)} = \delta_{rs} \delta^D(x - y)$), the chain rule, and integration by parts) that in this special case, the EOM (1) reduces to the Lagrange EOM that someone made you memorize:

$$0 = \frac{\partial \mathcal{L}}{\partial \phi_r} - \partial^\mu \frac{\partial \mathcal{L}}{\partial (\partial^\mu \phi_r)} \quad \text{for each } r.$$

(Note that since we are interested here in the bulk equations of motion, we ignore boundary terms unless we are interested in field theory on a space with boundary. That is a worthy subject but an unnecessary complication for now.)

3. **Momentum.** In this problem we consider a scalar field theory in d spatial dimensions. The scalar field is expanded in normal modes as usual:

$$\phi(\vec{x}) = \int d^d k \sqrt{\frac{\hbar}{2\omega_k}} \left(e^{i\vec{k}\cdot\vec{x}} a_k + e^{-i\vec{k}\cdot\vec{x}} a_k^\dagger \right),$$

where $\int d^d k \dots \equiv \int \frac{d^d k}{(2\pi)^d} \dots$. Consider the operator

$$\vec{P} \equiv \int d^d k \hbar \vec{k} a_k^\dagger a_k.$$

Since I wrote an integral for k , assume we are in the thermodynamic limit, and interpret the $\mathbf{a}$ s to be normalized as $[\mathbf{a}_k, \mathbf{a}_{k'}^\dagger] = (2\pi)^d \delta^d(k - k')$, a Dirac delta.

- (a) What is $\vec{P} |n_{k_1}, n_{k_2} \dots\rangle$ (where $|n_{k_1}, n_{k_2} \dots\rangle$ is a state with n_{k_i} phonons of wavenumber k_i)?
- (b) Find $[\vec{P}, a_k^\dagger]$, and $[\vec{P}, a_k]$.
- (c) Show, using [3b](#) and the mode expansion of the scalar field, that

$$[\vec{P}, \phi(x)] = i\hbar \vec{\nabla} \phi(x).$$

- (d) Conclude (using Taylor's theorem) that

$$e^{-i\vec{a}\cdot\vec{P}/\hbar} \phi(x) e^{i\vec{a}\cdot\vec{P}/\hbar} = \phi(x + \vec{a})$$

and that therefore $\vec{P}$ generates translations. Therefore $\vec{P}$ is the operator representing the momentum carried by the field (like the Poynting vector for the electromagnetic field).

- (e) Find $\vec{P} |\vec{k}_1, \vec{k}_2 \dots \vec{k}_n\rangle$, the action of this operator on a state of n phonons. Conclude that $\hbar \vec{k}$ is the momentum of the phonon labelled by wavenumber $\vec{k}$, as promised in lecture.
- (f) [Bonus problem] Show that $\vec{P}$ can be written as an integral of a local function of $\phi(x)$ and its derivatives.

4. Complex scalar field and antiparticles.

[This problem is related to Peskin problem 2.2.] So far we've discussed scalar field theory with one *real* scalar field. The particles created by such a field are their own antiparticles.

To understand this statement better, consider a scalar field theory in $d + 1$ dimensions with *two* real fields ϕ_1, ϕ_2 . Organize them into one complex field $\Phi \equiv \frac{1}{\sqrt{2}} (\phi_1 + \mathbf{i}\phi_2)$, with $\Phi^* = \frac{1}{\sqrt{2}} (\phi_1 - \mathbf{i}\phi_2)$, and let

$$S[\Phi, \Phi^*] = \int d^d x dt \left(\frac{1}{2} \mu \partial_t \Phi \partial_t \Phi^* - \frac{1}{2} \mu v^2 \vec{\nabla} \Phi \cdot \vec{\nabla} \Phi^* - V(\Phi^* \Phi) \right)$$

for some potential function $V(x)$.

(a) Show that

$$S[\Phi, \Phi^*] = \int \left(\sum_{i=1,2} \left(A (\partial_t \phi_i)^2 - B \vec{\nabla} \phi_i \cdot \vec{\nabla} \phi_i \right) - V((\phi_1^2 + \phi_2^2)/2) \right),$$

and where A, B are constants you must determine. If $V(q^2) = \frac{1}{2} \tilde{m}^2 q^2$, notice that the action is just the sum of two copies of the action of the theory we considered previously.

(b) Show by doing the Legendre transformation that the associated hamiltonian is

$$\mathbf{H} = \int d^d x \left(C \Pi \Pi^* + D \vec{\nabla} \Phi \cdot \vec{\nabla} \Phi^* + V(\Phi \Phi^*) \right)$$

where C, D are constants you must determine, and the canonical momenta are

$$\Pi = \frac{\partial \mathcal{L}}{\partial \dot{\Phi}} = \frac{1}{2} \mu \dot{\Phi}^*, \quad \Pi^* = \frac{\partial \mathcal{L}}{\partial \dot{\Phi}^*} = \frac{1}{2} \mu \dot{\Phi}$$

with the Lagrangian density $\mathcal{L}$ defined by $S = \int dt d^d x \mathcal{L}$.

(c) This theory has a continuous symmetry under which $\Phi \rightarrow e^{i\alpha} \Phi, \Phi^* \rightarrow e^{-i\alpha} \Phi^*$ with α a real constant. Show that the action S does not change if I make this replacement. ¹

(d) The existence of a continuous symmetry means a conserved charge – a hermitian operator that commutes with the Hamiltonian, which generates the symmetry (this is the quantum version of the Noether theorem). Show that

$$\mathbf{q} \equiv \int d^d x \mathbf{i} (\Phi^* \Pi^* - \Pi \Phi)$$

¹This is called a U(1) symmetry: it is a unitary rotation (hence ‘U’) on a one-dimensional (hence ‘(1)’) complex vector. Notice that on the real components ϕ_1, ϕ_2 it acts as a two-dimensional rotation:

$$\begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix} \rightarrow \begin{pmatrix} \cos \alpha & -\sin \alpha \\ \sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} \phi_1 \\ \phi_2 \end{pmatrix}.$$

The name for this group is SO(2). So U(1) is the same as SO(2).

generates this transformation, in the sense that the change in the field under a transformation with infinitesimal α is

$$\delta\Phi = \mathbf{i}\alpha\Phi = -\mathbf{i}\alpha[\mathbf{q}, \Phi], \quad \text{and} \quad \delta\Phi^* = -\mathbf{i}\alpha\Phi^* = -\mathbf{i}\alpha[\mathbf{q}, \Phi^*].$$

Show that $[\mathbf{q}, \mathbf{H}] = 0$.

- (e) For the case where $V(\Phi\Phi^*) = \frac{1}{2}\tilde{m}^2\Phi\Phi^*$ the hamiltonian is quadratic. Diagonalize it in terms of *two sets* of creation operators and annihilation operators. Work in the continuum. You should find something of the form

$$\Phi(x) = A \sum_k \frac{1}{\sqrt{\omega_k}} \left(e^{\mathbf{i}kx} \mathbf{a}_k + e^{-\mathbf{i}kx} \mathbf{b}_k^\dagger \right) \quad (2)$$

(and a similar expression for Π) for some constant A that you should determine (you may wish to do the next part first).

- (f) Write the canonical commutators

$$[\Phi(x), \Pi(x')] = \mathbf{i}\hbar\delta^d(x - x'), \quad [\Phi(x), \Pi^*(x')] = 0$$

(and the hermitian conjugate expressions) in terms of $\mathbf{a}$ and $\mathbf{b}$.

- (g) Rewrite $\mathbf{q}$ in terms of the mode operators.

- (h) Evaluate the charge of each type of particle created by $\mathbf{a}_k^\dagger$ and $\mathbf{b}_k^\dagger$ (*i.e.* find $[\mathbf{q}, \mathbf{a}^\dagger]$).

I claim that the particle created by $\mathbf{a}^\dagger$ is the antiparticle of that created by $\mathbf{b}^\dagger$ in the sense that they have opposite quantum numbers. This means that we can add terms to the hamiltonian by which they can annihilate each other, without breaking any symmetries. What might such a term look like? Write something local in position space, *i.e.* of the form $\Delta H = \int d^d x V(\Phi, \Phi^\dagger)$.